
Quantum mechanics II, Problems 10 : Characters and Lie Algebra Basics

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Problem 1 : Irreps of C_{3v}

The aim of this exercise is to consider two representations of the group C_{3v} . We will start by finding the representation R of the group on the vector space \mathbb{R}^2 . Then we will find the representation P_R of the group on the function space generated by $\Psi_1(\mathbf{r}) = x^2e^{-r}$, $\Psi_2(\mathbf{r}) = y^2e^{-r}$, $\Psi_3(\mathbf{r}) = 2xye^{-r}$. We will show that the representation P_R is reducible. We will establish the connection with the representation R of dimension 2.

1. Consider the vector space \mathbb{R}^2 with vectors (x, y) . Derive the representation of $R(\sigma_1)$ and $R(C_3)$ in this space. Then deduce the group multiplication table to find $R(u)$, $\forall u \in C_{3v}$. We will assume that this representation is unitary and irreducible, which can be demonstrated by Schur's theorem.
2. Consider now the vector space of functions \mathcal{H} , generated by functions :

$$\begin{aligned}\Psi_1(\mathbf{r}) &= x^2e^{-r} \\ \Psi_2(\mathbf{r}) &= y^2e^{-r} \\ \Psi_3(\mathbf{r}) &= 2xye^{-r}\end{aligned}$$

where $r = |\mathbf{r}| = \sqrt{x^2 + y^2}$, with the scalar product :

$$\langle \Psi_\alpha | \Psi_\beta \rangle = \int d^2\mathbf{r} \Psi_\alpha^*(\mathbf{r}) \Psi_\beta(\mathbf{r}).$$

Written as matrices, the group representation C_{3v} is defined as follows :

$$P_{R(u)}\Psi(\mathbf{r}) \equiv \Psi(R^{-1}(u)\mathbf{r}), \forall u \in C_{3v},$$

where $R(u)$ are the matrices derived in point (a) (in quantum mechanics, for example, the wave function of a particle obeys this transformation law following a rotation of the reference frame). Show that it is a representation of the group, and that its matrices are not all unitary.

3. Show that the representation $P_{R(u)}$ is reducible by identifying an invariant subspace.
4. Hence show that the representation $P_{R(u)}$ can be written as a direct sum of a 2D and 1D irreducible representations.
5. Construct the character table of C_{3v} !

Problem 2 : Lie-Algebras and Infinitesimal Generators

This problem is intended to get you familiar with the basics of Lie Algebras and help you understand the relationship between $SU(2)$ and $SO(3)$ (which in turn will help you (hopefully!) have a better understanding of why the Bloch sphere representation of quantum states works).

1. Compute a 3D representation of the basis of the Lie-Algebra of $SO(3)$. Then compute the structure constants (commutator) among the basis elements. Show that the representation $\rho : SO(3) \rightarrow GL(\mathbb{R}^3)$ with $\rho(A) = e^{a_i X_i}$, $A \in SO(3)$, $a_i \in \mathbb{R}$ representing the rotation (for example angles) and X_i the basis elements of $\mathfrak{so}(3)$, is a valid representation of $SO(3)$ (this is called the fundamental representation of $SO(3)$).

Hint : The Lie-Algebra is formally defined as the tangent space to the Lie-Group at the identity element. In practice we can use this to compute the Lie-Algebra by means of the exponential map. In fact any element $A \in G$ can be written as $A(t) = e^{tX}$, $X \in \mathfrak{g}$, where \mathfrak{g} is the Lie-Algebra. Therefore one can access elements by looking at : $\frac{d}{dt}A(t)|_{t=0} \in \mathfrak{g}$.

2. Do the same for $SU(2)$.
3. It can be shown that the finite dimensional, irreducible representations of $SO(3)$ all have odd dimensions. They can be constructed with the help of the well known ladder operators $L_{\pm} = L_x \pm iL_y$ where $L_{x,y,z}$ form a basis of $\mathfrak{so}(3)$. In the basis in which L_z is diagonal i.e. $\{|l, m\rangle\}$, $L_z |l, m\rangle = m |l, m\rangle$, it can be shown that $L_{\pm} |l, m\rangle = \sqrt{(l+1 \pm m)(l \mp m)} |l, m \pm 1\rangle$. Use this to i) compute the 3D irreducible representation of $SO(3)$ and ii) the 5D irreducible representation of $SO(3)$.

Hint : Compute the matrix representation of the ladder operators L_{\pm} in the basis $\{|l, m\rangle\}$ and construct $L_{x,y}$ from there.

4. For $SU(2)$ the irreducible representations can have any dimension and we have the same ladder operators as for $SO(3)$. i) Compute the 2D irreducible representation of $SU(2)$. ii) Compute the 3D irreducible representation of $SU(2)$. Compare with the one from $SO(3)$.
5. Is the 3D representation of $SO(3)$ you have derived also a representation of $SU(2)$? Explain why this makes sense with respect to the Bloch sphere.

(Bonus - non examinable - are all representations of $SU(2)$ also representations of $SO(3)$? What about vice versa?)